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## A RIEMANN RESTRICTED CHARACTERIZATION OF THE QUASILINEARITY HIERARCHY: SOME REMARKS

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*To Professor Dan Pascali, at his 70's anniversary*

### Abstract

A Riemann restricted version of the details connected with Lax's genuine nonlinearity / linear degeneracy ([3]) is considered. An example regarding the interface between weak quasilinearity and linearity is included. A case is presented for which the weak quasilinearity *degenerates* into a strong quasilinearity.

### Introduction. A quasilinearity hierarchy

We consider the one-dimensional strictly hyperbolic system

$$\frac{\partial u}{\partial t} + a(u) \frac{\partial u}{\partial x} = 0, \quad (1)$$

and let  $\overset{i}{R}(u)$ ,  $\overset{i}{L}(u)$ , and  $\lambda_i(u)$ , of indices  $i = 1, \dots, n$ , be, respectively, the right eigenvectors, the left eigenvectors, and the eigenvalues of the matrix  $a(u)$ .

TERMINOLOGY 1 ([3]). For a strictly hyperbolic system (1) we say that an index  $i$  is *genuinely nonlinear* in  $\mathcal{R} \subset H$  ( $H$  is the hodograph space) if for it

$$\overset{i}{R}(u) \cdot \text{grad}_u \lambda_i(u) \neq 0, \quad u \in \mathcal{R} \quad (2)$$

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and it is *linearly degenerated* in  $\mathcal{R} \subset H$  if for it

$$\overset{i}{R}(u) \cdot \text{grad}_u \lambda_i(u) \equiv 0, \quad u \in \mathcal{R} \quad (3)$$

TERMINOLOGY 2. The quasilinearity associated to (1) is said to be *strong* if all the indices associated to (1) are genuinely nonlinear, *medium* if only a part of the indices associated to (1) are genuinely nonlinear [the other being linearly degenerate], and *weak* if all the indices associated to (1) are linearly degenerate.

## Riemann Invariants

Let us consider the  $(n - 1)$ -dimensional hypersurface

$$v_k(u) = \text{constant} = v_k(u^*) \quad (4)$$

through  $u^* \in \mathcal{R}$  at the points of which the normal direction is given by the left eigenvector  $\overset{k}{L}(u)$ :

$$\frac{\partial v_k}{\partial u_l} = \alpha_k(u) \overset{k}{L}_l(u) \quad 1 \leq l \leq n. \quad (5)$$

It results from (5) that the reality of a characteristic hypersurface (4) depends on the integrable character of the Pfaff form

$$\sum_{l=1}^n \overset{k}{L}_l(u) du_l. \quad (6)$$

In terms of (6), the geometrical restrictions mentioned here above appear to be integrability restrictions. A hypersurface (4), (5) is called a *Riemann invariant* of (1).

We use (5) to compute for each  $k$ :

$$\begin{aligned} \alpha_k(u) \overset{k}{L}(u) \left[ \frac{\partial u}{\partial t} + a(u) \frac{\partial u}{\partial x} \right] &= \sum_{j=1}^n \frac{\partial v_k}{\partial u_j} \cdot \frac{\partial u_j}{\partial t} + \lambda_k(u) \alpha_k(u) \overset{k}{L}(u) \frac{\partial u}{\partial x} \\ &= \frac{\partial v_k}{\partial t} + \bar{\lambda}_k(v) \frac{\partial v_k}{\partial x}, \quad 1 \leq k \leq n \end{aligned}$$

and notice that the Riemann invariants  $v_1, \dots, v_n$  satisfy a *diagonal* system

$$\frac{\partial v_k}{\partial t} + \bar{\lambda}_k(v) \frac{\partial v_k}{\partial x} = 0, \quad 1 \leq k \leq n; \quad \bar{\lambda}_k(v) \equiv \lambda_k[u(v)] \quad (7)$$

associated to (1). In fact, a Riemann invariant  $v_k$  is constant, cf. (7), on each characteristic line of index  $k$ . In (7) we used, in order to define  $\bar{\lambda}_k$ , the non-singular transformation character of the connection between  $v$  and  $u$  around each point of  $\mathcal{R}$ .

### Riemann restricted genuine nonlinearity / linear degeneracy

We also compute from (7)

$$\frac{\partial \lambda_i}{\partial u_j} = \sum_{k=1}^n \frac{\partial \bar{\lambda}_i}{\partial v_k} \cdot \frac{\partial v_k}{\partial u_j} = \sum_{k=1}^n \alpha_k(u) L_j^k(u) \frac{\partial \bar{\lambda}_i}{\partial v_k}$$

$$R^i(u) \cdot \text{grad}_u \lambda_i(u) = \sum_{j=1}^n R_j^i(u) \sum_{k=1}^n \alpha_k(u) L_j^k(u) \frac{\partial \bar{\lambda}_i}{\partial v_k} = \bar{\alpha}_i(v) \frac{\partial \bar{\lambda}_i}{\partial v_i}$$

and respectively transcribe the restrictions of genuinely nonlinearity / linear degeneracy of an index  $i$  [see (2)/(3)] by

$$\frac{\partial \bar{\lambda}_i}{\partial v_i} \neq 0, \quad v \text{ in } \bar{\mathcal{R}} \quad (8)$$

or

$$\frac{\partial \bar{\lambda}_i}{\partial v_i} \equiv 0, \quad v \text{ in } \bar{\mathcal{R}}. \quad (9)$$

### Some remarks

At this point we shall use (8), (9) in order to characterize the quasilinearity hierarchy (Terminology 2). A complete system of Riemann invariants always exists as  $n = 2$  and we notice that a representative and most suggestive characterization of the mentioned hierarchy can be done for  $n = 2$ . For  $n = 2$ , a *strong* quasilinearity means

$$\frac{\partial \bar{\lambda}_1}{\partial v_1} \neq 0, \quad \frac{\partial \bar{\lambda}_2}{\partial v_2} \neq 0 \quad \text{in } \bar{\mathcal{R}}$$

a *medium* quasilinearity requires

$$\frac{\partial \bar{\lambda}_1}{\partial v_1} \neq 0, \quad \frac{\partial \bar{\lambda}_2}{\partial v_2} \equiv 0 \quad \text{or} \quad \frac{\partial \bar{\lambda}_1}{\partial v_1} \equiv 0, \quad \frac{\partial \bar{\lambda}_2}{\partial v_2} \neq 0 \quad \text{in } \bar{\mathcal{R}}$$

and a *weak* quasilinearity has the signification

$$\frac{\partial \bar{\lambda}_1}{\partial v_1} \equiv 0, \quad \frac{\partial \bar{\lambda}_2}{\partial v_2} \equiv 0 \quad \text{in } \bar{\mathcal{R}}. \quad (10)$$

A nontrivial form of (10) is complementarily characterized by

$$\frac{\partial \bar{\lambda}_1}{\partial v_2} \neq 0, \quad \frac{\partial \bar{\lambda}_2}{\partial v_1} \neq 0 \quad \text{in } \bar{\mathcal{R}}. \quad (11)$$

As (10) and (11) hold we set

$$r = \bar{\lambda}_2(v_1), \quad s = \bar{\lambda}_1(v_2)$$

in order to transform the corresponding system (7) into

$$\frac{\partial r}{\partial t} + s \frac{\partial r}{\partial x} = 0, \quad \frac{\partial s}{\partial t} + r \frac{\partial s}{\partial x} = 0. \quad (12)$$

Then, we calculate from (12)

$$\frac{\partial r}{\partial t} \frac{\partial s}{\partial x} - \frac{\partial r}{\partial x} \frac{\partial s}{\partial t} = (r - s) \frac{\partial r}{\partial x} \frac{\partial s}{\partial x} \quad (13)$$

and

$$\frac{\partial r}{\partial t} + r \frac{\partial r}{\partial x} = (r - s) \frac{\partial r}{\partial x}, \quad \frac{\partial s}{\partial t} + s \frac{\partial s}{\partial x} = (s - r) \frac{\partial s}{\partial x}. \quad (14)$$

If we weaken the restriction (11), allowing for example that

$$\frac{\partial \bar{\lambda}_1}{\partial v_2} \equiv 0 \quad \left[ \text{yet } \frac{\partial \bar{\lambda}_2}{\partial v_2} \neq 0 \right] \quad \text{in } \bar{\mathcal{R}},$$

then (13) takes the form

$$\frac{\partial r}{\partial t} + h \frac{\partial r}{\partial x} = 0, \quad \frac{\partial s}{\partial t} + r \frac{\partial s}{\partial x} = 0, \quad \text{constant } h$$

which reduces to a *linear* equation:

$$\frac{\partial s}{\partial t} + r_0(x - ht) \frac{\partial s}{\partial x} = 0.$$

It appears that the restrictions (10) and (11) characterize the *lowest* level of nonlinearity in the weak quasilinearity connected with  $n = 2$ .

We finally notice that a solution of (12) for which  $r \equiv s$  corresponds to a *degeneration* of the weakly quasilinear system (12). In this degeneration the two equations (12) become coincident in the genuinely nonlinear equation

$$\frac{\partial r}{\partial t} + r \frac{\partial r}{\partial x} = 0.$$

The mentioned degeneration implies a replacement of a  $n = 2$  linear degeneracy by a  $n = 1$  genuine nonlinearity.

## References

- [1] L.F. Dinu, Some remarks concerning the Riemann invariance. Burnat–Peradzynski and Martin approaches, *Revue Roumaine Math. Pures Appl.*, **35**(1990), 203–234.
- [2] L.F. Dinu, *Mathematical concepts in nonlinear gas dynamics*, CRC Press [to appear].
- [3] P.D. Lax, Hyperbolic systems of conservation laws (II), *Comm. Pure and Appl. Math.*, **10**(1957), 537–566.
- [4] R. von Mises, *Mathematical theory of compressible fluid flow*, completed by Hilda Geiringer and G.S.S. Ludford, Academic Press, New York, 1958.
- [5] G.B. Riemann, Über die Fortpflanzung ebener Luftwellen von endlicher Schwingungsweite, *Ahandl. Kön. Ges. Wiss. Göttingen, Math.–physik. Kl.*, **8**(1858/9), 43–65.
- [6] J.A. Smoller, *Shock waves and reaction-diffusion equations*, Springer Verlag, Berlin – Heidelberg – New York, 1983.

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