



# On geometric structures associated with triple systems

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## Abstract

In this paper, we show that geometric phenomena can be characterized using the concept of triple systems. In particular, we study a complex structure associated with triple systems.

## 1 Introduction

In future, it is expected that triple systems will be useful for the characterization of geometric structures in mathematics and physics as well as that of (classical) Yang-Baxter equations ([10],[19],[22]).

Our aim is to use triple systems to investigate the characterization of differential geometry and mathematical physics from the viewpoint of non-associative algebras that contain a class of Lie algebras or Jordan algebras ([6],[7],[8],[13],[16],[17],[18]). In this paper, we will study about the Lie super-algebras or Lie algebras associated with triple systems, and the relationship with several vector subspaces associated with these algebras. We show that these subspaces have a complex structure as well as differential geometry, and we provide an algebraic characterization of their subspaces. Furthermore, for  $B_3$ -type Lie algebras, we will give some examples of triple systems and their correspondence with extended Dynkin diagrams.

A  $(2\nu+1)$ -graded Lie algebra is a Lie algebra of the form  $g = \bigoplus_{k=-\nu}^{\nu} g_k$  such that  $[g_k, g_l] \subset g_{k+l}$ . It is well known that 3-graded Lie algebras are essentially bijective with certain theoretic objects called Jordan triple systems or Jordan

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pairs. Kantor remarked that more general graded Lie algebras correspond to generalized Jordan triple systems. In particular, the graded Lie algebra  $g$  (or  $L$ ) given by

$$g_{-2} \oplus g_{-1} \oplus g_0 \oplus g_1 \oplus g_2 \quad (\text{or } L_{-2} \oplus L_{-1} \oplus L_0 \oplus L_1 \oplus L_2)$$

has the structure of a triple product on the subspace  $g_{-1}$ , and is known as a generalized Jordan triple system (GJTS) of the second order or a (-1,1)-Freudenthal-Kantor triple system (F-K.t.s.)([8],[14]). Also  $g_{-1} \oplus g_1$  has the structure of a Lie triple system (in particular, the system over a real number is known to correspond with a symmetric Riemannian space by means of a totally geodesic manifold notation). We will discuss the corresponding geometrical object by using these triple systems. The notation and terminology used for the geometry can be found in ([4], [5], [21]). We will often use the symbols  $g$  and  $L$  to denote a Lie algebra or Lie superalgebra as is conventionally used ([2],[3],[7],[25]).

From the viewpoint of an algebraic study, our purpose is to propose a unified structural theory for triple systems in nonassociative algebras. In previous works ([12],[14],[15]), we have studied the Peirce decomposition of a GJTS  $U$  of the second order by employing a tripotent element  $e$  of  $U$  (for a tripotent element,  $\{eee\} = e$ ).

The Peirce decomposition of  $U$  is described as follows:

$$U = U_{00} \oplus U_{\frac{1}{2}\frac{1}{2}} \oplus U_{11} \oplus U_{\frac{3}{2}\frac{3}{2}} \oplus U_{-\frac{1}{2}0} \oplus U_{01} \oplus U_{\frac{1}{2}2} \oplus U_{13},$$

where  $L(a) = \{eea\} = \lambda a$  and  $R(a) = \{aee\} = \mu a$  if  $a \in U_{\lambda\mu}$ .

These viewpoints as above have formed the basis of our study on triple systems. However in this note, our consideration will mainly be from a geometrical viewpoint.

We are concerned with triple systems and algebras which have finite dimensionality over a field  $\Phi$  of characteristic  $\neq 2$  or  $3$ , unless otherwise specified.

## 2 Definitions and Preamble

To make this paper as self-contained as possible, we first recall the definition of a generalized Jordan triple system of the second order (hereafter, referred to as a GJTS of 2nd order), and the construction of Lie algebras associated with a GJTS of 2nd order.

A vector space  $V$  over a field  $\Phi$ , endowed with a trilinear operation  $V \times V \times V \rightarrow V$ ,  $(x, y, z) \mapsto \{xyz\}$ , is said to be a *GJTS of 2nd order* if the following two conditions are satisfied:

- (J1)  $\{ab\{xyz\}\} = \{\{abx\}yz\} - \{x\{bay\}z\} + \{xy\{abz\}\}$  (GJTS)
- (K1)  $K(K(a, b)x, y) - L(y, x)K(a, b) - K(a, b)L(x, y) = 0$  (2nd order),

where  $L(a, b)c = \{abc\}$  and  $K(a, b)c = \{acb\} - \{bca\}$ .

**Remark.** If  $K(a, b) \equiv 0$  (identically zero), then this triple system is a Jordan triple system (JTS), i.e., it satisfies the relations  $\{acb\} = \{cba\}$  and *GJTS*.

We can also generalize the concept of the GJTS of 2nd order as follows (for examples, see [8],[9],[11],[16] and references therein).

For  $\varepsilon = \pm 1$  and  $\delta = \pm 1$ , if the triple product satisfies

$$\begin{aligned} (abxyz) &= ((abx)yz) + \varepsilon(x(bay)z) + (xy(abz)), \\ K(K(a, b)c, d) - L(d, c)K(a, b) + \varepsilon K(a, b)L(c, d) &= 0, \end{aligned}$$

where  $L(x, y)z = (xyz)$  and  $K(a, b)c = (acb) - \delta(bca)$ , then it is said to be an  $(\varepsilon, \delta)$ -Freudenthal-Kantor triple system (hereafter abbreviated as an  $(\varepsilon, \delta)$ -F-K.t.s).

If  $Id \in \mathbf{k} = \{K(a, b)\}_{span}$ , it is said to be *unitary*.

Furthermore, if the  $(\varepsilon, \delta)$ -F-K.t.s satisfies

$$dim_{\Phi}\{K(a, b)\}_{span} = dim_{\Phi}\{< a, b >\}_{span} = 1,$$

where  $< a, b > \in \Phi^*$ , then it is said to be *balanced*.

**Remark.** We set  $S(x, y) := L(x, y) + \varepsilon L(y, x)$ , and  $A(x, y) := L(x, y) - \varepsilon L(y, x)$ , then this  $S(x, y)$  (resp.  $A(x, y)$ ) is a derivation (resp. anti-derivation) of  $U(\varepsilon, \delta)$ .

**Remark.** Following the notation of the metasymplectic geometry due to H.Freudenthal, our concept means that

$$PxQ \dots \dots \text{derivation } S(x, y) \text{ of } U(\varepsilon, \delta),$$

$\{P, Q\} \dots \dots \text{anti derivation } A(x, y) (= K(x, y)) \text{ of } U(\varepsilon, \delta) \dots \text{balanced type.}$

We generally denote the triple products by  $\{xyz\}$ ,  $(xyz)$ ,  $[xyz]$ , and  $< xyz >$ . Bilinear forms are denoted by  $< x|y >$ ,  $< x, y >$ , and  $B(x, y)$ .

**Remark.** Note that the concept of a GJTS of 2nd order coincides with that of a  $(-1, 1)$ -F-K.t.s. Thus we can construct simple Lie algebras ( $\delta = 1$ ) or superalgebras ( $\delta = -1$ ) from these triple systems with  $\delta = \pm 1$ , by means of the standard embedding method (for example, [2], [3], [8]-[12], [16], [17], [18], [23]).

**Proposition 1** ([9],[17]). *Let  $U(\varepsilon, \delta)$  be an  $(\varepsilon, \delta)$ -F-K.t.s. If  $J$  is an endomorphism of  $U(\varepsilon, \delta)$  such that  $J < xyz > = < JxJyJz >$  and  $J^2 = -\varepsilon\delta Id$ , then  $(U(\varepsilon, \delta), [xyz])$  is a Lie triple system (the case of  $\delta = 1$ ) or an anti-Lie triple system (the case of  $\delta = -1$ ) with respect to the product*

$$[xyz] := < xJyz > - \delta < yJxz > + \delta < xJzy > - < yJzx > .$$

**Corollary.** *Let  $U(\varepsilon, \delta)$  be an  $(\varepsilon, \delta)$ -F.K.t.s. Then the vector space  $T(\varepsilon, \delta) = U(\varepsilon, \delta) \oplus U(\varepsilon, \delta)$  becomes a Lie triple system (the case of  $\delta = 1$ ) or an anti-Lie triple system (the case of  $\delta = -1$ ) with respect to the triple product defined by*

$$\left[ \begin{pmatrix} a \\ b \end{pmatrix} \begin{pmatrix} c \\ d \end{pmatrix} \begin{pmatrix} e \\ f \end{pmatrix} \right] = \begin{pmatrix} L(a, d) - \delta L(c, b) & \delta K(a, c) \\ -\varepsilon K(b, d) & \varepsilon(L(d, a) - \delta L(b, c)) \end{pmatrix} \begin{pmatrix} e \\ f \end{pmatrix}.$$

Thus we can obtain the standard embedding Lie algebra (the case of  $\delta = 1$ ) or Lie superalgebra (the case of  $\delta = -1$ ),  $L(\varepsilon, \delta) = D(T(\varepsilon, \delta), T(\varepsilon, \delta)) \oplus T(\varepsilon, \delta)$ , associated with  $T(\varepsilon, \delta)$ , where  $D(T(\varepsilon, \delta), T(\varepsilon, \delta))$  is the set of inner derivations of  $T(\varepsilon, \delta)$ . That is, these vector spaces  $D(T(\varepsilon, \delta), T(\varepsilon, \delta))$  and  $T(\varepsilon, \delta)$  imply

$$D(T(\varepsilon, \delta), T(\varepsilon, \delta)) := \begin{pmatrix} L(a, b) & \delta K(c, d) \\ -\varepsilon K(e, f) & \varepsilon L(b, a) \end{pmatrix}_{span} \quad \text{and}$$

$$T(\varepsilon, \delta) := \left\{ \begin{pmatrix} x \\ y \end{pmatrix} \mid x, y \in U(\varepsilon, \delta) \right\}_{span}, \quad (\text{denoted by } T(U)).$$

In fact, we have

$$L_0 = \left\{ \begin{pmatrix} L(a, b) & 0 \\ 0 & \varepsilon L(b, a) \end{pmatrix} \right\}_{span} = \{L(a, b)\}_{span},$$

$$L_{-2} = \left\{ \begin{pmatrix} 0 & \delta K(c, d) \\ 0 & 0 \end{pmatrix} \right\}_{span} = \{K(c, d)\}_{span} \text{ and } L_0 = DerU \oplus Anti - DerU.$$

**Remark.** For the standard embedding algebras obtained from these triple systems, note that

$$L(\varepsilon, \delta) := L_{-2} \oplus L_{-1} \oplus L_0 \oplus L_{-1} \oplus L_{-2}$$

(or  $g = g_{-2} \oplus g_{-1} \oplus g_0 \oplus g_1 \oplus g_2$ ) is a 5-graded Lie algebra or Lie superalgebra such that

$$L_{-1} = g_{-1} = U(\varepsilon, \delta)$$

and

$$DerT(U) := D(T(\varepsilon, \delta), T(\varepsilon, \delta)) = L_{-2} \oplus L_0 \oplus L_{-2}$$

with

$$[L_i, L_j] \subseteq L_{i+j}.$$

Also, denote  $L(\varepsilon, \delta)$  by  $L(U)$ .

For the correspondence of the (1,1) F.K.t.s with the (-1,1) F.K.t.s, we obtain the following.

**Proposition 2.** *Let  $(U, \langle xyz \rangle)$  be a  $(1, 1)$  F-K.t.s. If there is an endomorphism  $P$  of  $U$  such that  $P \langle xyz \rangle = \langle PxPyPz \rangle$  and  $P^2 = -Id$ , then  $(U, \{xyz\})$  is a GJTS of 2nd order (that is,  $(-1, 1)$ -F-K.t.s.) with respect to the new product defined by  $\{xyz\} := \langle xPyz \rangle$ .*

*Proof.* Suppose that the product  $\langle xyz \rangle$  is satisfying the two relations;

$$\langle ab \langle xyz \rangle \rangle = \langle \langle abx \rangle yz \rangle + \langle x \langle bay \rangle z \rangle + \langle xy \langle abz \rangle \rangle,$$

$$K(K(a, b)c, d) - L(d, c)K(a, b) + K(a, b)L(c, d) = 0,$$

where  $L(a, b)c = \langle abc \rangle$  and  $K(a, b)c = \langle acb \rangle - \langle bca \rangle$ .

By means of  $P^2 = -Id$  and the definition  $\{xyz\} = \langle xPyz \rangle$ , it is enough to show that the triple product  $\{xyz\}$  is satisfied the two identities;

$$\{ab\{xyz\}\} = \{\{abx\}yz\} - \{x\{bay\}z\} + \{xy\{abz\}\},$$

$$K(K(a, b)x, y) - L(y, x)K(a, b) - K(a, b)L(x, y) = 0$$

where  $L(x, y)z = \{xyz\}$  and  $K(a, b)c = \{acb\} - \{bca\}$ .

Thus, these relations is obtained by straightforward calculatins, This completes the proof.

We now give an explicit example of a JTS and a Lie triple system.

**Example.** Let  $U$  be a vector space with a symmetric bilinear form  $\langle , \rangle$ . Then the triple system  $(U, [xyz])$  is a Lie triple system with respect to the product

$$[xyz] = \langle y, z \rangle x - \langle z, x \rangle y.$$

That is, this triple system is induced from the JTS

$$\{xyz\} = \frac{1}{2}(\langle x, y \rangle z + \langle y, z \rangle x - \langle z, x \rangle y)$$

by means of

$$[xyz] = \{xyz\} - \{yxz\}.$$

### 3 A complex structure associated with triple systems

In this section, we will discuss with a complex structure on the following vector space;

$$T(\varepsilon, \delta) = g_{-1} \oplus g_1.$$

We set

$$E = \begin{pmatrix} 0 & Id \\ 0 & 0 \end{pmatrix}, F = \begin{pmatrix} 0 & 0 \\ Id & 0 \end{pmatrix}, H = \begin{pmatrix} Id & 0 \\ 0 & -Id \end{pmatrix}, J = \delta E - \varepsilon F.$$

Then we obtain by straightfoward calculations,

$$H = [E, F], [H, E] = 2E, [H, F] = -2F, J^2 = -\delta\varepsilon \begin{pmatrix} Id & 0 \\ 0 & Id \end{pmatrix},$$

where  $J$  is an operator on  $T(\varepsilon, \delta)$ .

Next we define the Nijenhuis operator on  $T(\varepsilon, \delta)$  as below.

$$N(X, Y) = [JX, JY] - J[JX, Y] - J[X, JY] + J^2[X, Y], \quad X, Y \in T(\varepsilon, \delta).$$

We study the cases of  $\varepsilon\delta = 1$ , which can be considered to have an almost complex structure, since  $J^2 = -Id$ . However, we do not deal with the case of  $\varepsilon\delta = -1$  called a paracomplex structure, which will be considered in future work.

By a long but straightfoward calculations, proof omitted, we obtain the following Proposition.

**Proposition 3.** *Let  $U$  be an  $(\varepsilon, \delta)$ -F-K.t.s. and its operations be given by the above definitions. The following are equivalent:*

- (i)  $N(X, Y) = 0$ ,
- (ii)  $L(y, x) - \delta L(x, y) = K(x, y)$ .

From these results as well as the differential geometry, we conclude that there exists a complex structure on  $T(\varepsilon, \delta)$  if  $L(y, x) - \delta L(x, y) = K(x, y)$ .

Following [16], we exhibit examples of a  $(-1, -1)$ -F-K.t.s with a complex structure, is known as an antistructurable algebra.

**Remark.** By the well known fact that the Lie triple systems have a correspondence with symmetric spaces,  $T(\varepsilon, \delta)$  is closely related to symmetric space.

By using above  $J$ , we may define an operator  $\tilde{J}$  on  $L(\varepsilon, \delta)$  as follows:

$$\begin{pmatrix} L(a, b) & \delta K(c, d) \\ -\varepsilon K(e, f) & \varepsilon L(b, a) \end{pmatrix} \oplus \begin{pmatrix} x \\ y \end{pmatrix} \longrightarrow J \begin{pmatrix} L(a, b) & \delta K(c, d) \\ -\varepsilon K(e, f) & \varepsilon L(b, a) \end{pmatrix} J^{-1} \oplus J \begin{pmatrix} x \\ y \end{pmatrix}.$$

Then this  $\tilde{J}$  satisfies

$$\tilde{J}[X, Y] = [\tilde{J}X, \tilde{J}Y],$$

for all  $X, Y \in L(\varepsilon, \delta)$ .

Thus  $\tilde{J}$  is said to be an automorphism of  $L(\varepsilon, \delta)$  induced from the almost complex structure  $J$ .

**Remark.** Note that if  $U$  is unitary, then  $L(\varepsilon, \delta)$  contains the subalgebra  $sl_2 = \{H, E, F\}_{span}$ , because  $Id \in \mathfrak{k} = \{K(a, b)\}_{span} = g_{-2}$ .

In this section's final comment, we will introduce a deformed Nijenhuis operator on  $H(\varepsilon, \delta) = g_0 \oplus g_1 \oplus g_2$ .

$$N_H(X, Y) = [J_H X, J_H Y] - J_H [J_H X, Y] - J_H [X, J_H Y] + J_H^2 [X, Y],$$

$X, Y \in H(\varepsilon, \delta)$ , where  $J_H : (x_0, x_1, x_2) \rightarrow (\alpha x_0, \alpha x_1, \alpha x_2)$ .

**Proposition 4.** *For  $H(\varepsilon, \delta)$ , if there is a element  $\alpha \in \Phi$  such that  $\alpha^2 = -1$ , then the space  $H(\varepsilon, \delta)$  has a complex structure.*

*Proof.* From  $\alpha^2 = -1$ , it follows that the identity  $N_H(X, Y) = 0$  holds. Hence, we have a complex structure.

#### 4 Several algebraic structures

In this section, we will study some algebraic structures on (i)  $T(\varepsilon, \delta)$ , (ii)  $B(\varepsilon, \delta) = g_{-2} \oplus g_{-1} \oplus g_1 \oplus g_2$  and (iii)  $H(\varepsilon, \delta)$ .

We denote any element of  $g_{-2} \oplus g_{-1} \oplus g_0 \oplus g_1 \oplus g_2$  by  $(x_{-2}, x_{-1}, x_0, x_1, x_2)$  or  $x_{-2} + x_{-1} + x_0 + x_1 + x_2$ .

Before going into further details, we recall the definition of a generalized structurable algebra  $A$  (e.g., [13]) as having a bilinear derivation  $D(x, y)$  satisfying

$$\begin{aligned} D(x, y)(uv) &= (D(x, y)u)v + u(D(x, y)v) \\ D(xy, z) + D(yz, x) + D(zx, y) &= 0. \end{aligned}$$

**Example.([13])** Any Lie algebra, Jordan algebra, alternative algebra or structurable algebra is a generalized structurable algebra.

In fact, the derivation in the case of a Lie algebra is  $ad[x, y]$  and in the case of a Jordan algebra is  $[L(x), R(y)]$  respectively.

By means of the concept of a generalized structurable algebra, we introduce the following algebraic structure.

For (i)  $T(\varepsilon, \delta)$ , by defining

$$X \circ Y = 0 \quad \text{and} \quad D(X, Y)Z = [XYZ],$$

i.e.,  $D(X, Y) = ad([x_{-1}, y_1] + [x_1, y_{-1}])$ , where  $X = x_{-1} + x_1$  and  $Y = y_{-1} + y_1$ , we obtain a structure with a trivial algebra which bilinear product  $\circ$  is identically zero. However, this space is a Lie triple system ( $\delta = 1$ ) or an anti-Lie triple system ( $\delta = -1$ ) (c.f., section 2), that is,

$$\begin{aligned} D(X, Y)Z &= [XYZ] = -\delta[YXZ] = -\delta D(Y, X)Z, \\ D(X, Y)Z + D(Y, Z)X + D(Z, X)Y &= 0, \\ [D(X, Y), D(U, V)] &= D(D(X, Y)U, V) + D(U, (D(X, Y)V)). \end{aligned}$$

**Remark.** Note that

$$D(X, Y) = -\delta D(Y, X) \quad \text{for all } X, Y \in T(\varepsilon, \delta).$$

Next, for the case (ii)  $B(\varepsilon, \delta)$ , we define as below

$$X \circ Y = [x_{-1}, y_{-1}] + [x_{-1}, y_2] + [x_{-2}, y_1] + [x_1, y_{-2}] + [x_1, y_1] + [x_2, y_{-1}]$$

and

$$D_B(X, Y) = ad([x_{-1}, y_1] + [x_{-2}, y_2] + [x_1, y_{-1}] + [x_2, y_{-2}]),$$

where  $X = x_{-2} + x_{-1} + x_1 + x_2$ , and  $Y = y_{-2} + y_{-1} + y_1 + y_2$ .

This space  $B(\varepsilon, \delta)$  is a generalized structurable algebra with respect to the above product  $X \circ Y$  and derivation  $D_B(X, Y)$ .

**Remark.** Note that

$$D_B(X, Y) = \begin{cases} -\delta D_B(Y, X) & \text{if } X, Y \in g_{-1} \oplus g_1 \\ -D_B(Y, X) & \text{if } X \text{ or } Y \in g_{-2} \oplus g_2. \end{cases}$$

For (iii)  $H(\varepsilon, \delta)$ , we define as below

$$X \circ Y = [x_0, y_0] + [x_1, y_0] + [x_2, y_0] + [x_1, y_1] + [x_0, y_2] + [x_0, y_1]$$

$$D_H(X, Y) = ad([x_0, y_0]),$$

where  $X = x_0 + x_1 + x_2$  and  $Y = y_0 + y_1 + y_2$ .

This space  $H(\varepsilon, \delta)$  is a generalized structurable algebra with respect to the above product  $X \circ Y$  and the derivation  $D_H(X, Y) = ad([x_0, y_0])$ .

**Remark.** Note that

$$D_H(X, Y) = \begin{cases} -\delta D_H(Y, X) & \text{if } X, Y \in g_1 \\ -D_H(Y, X) & \text{if } X \text{ or } Y \in g_0 \oplus g_2. \end{cases}$$

## 5 Construction of $B_3$ -type Lie algebras from several triple systems

In this section, we will discuss the construction of simple  $B_3$ -type Lie algebras associated with several triple systems (the details will be described in a forthcoming paper).

- a) The case of a JTS,
- b) The case of a balanced GJTS,
- c) The case of a GJTS of 2nd order,
- d) The case of a derivation induced from a JTS.

To consider these cases, we start with an extended Dynkin diagram for a  $B_3$ -type Lie algebra.

$$\begin{array}{ccccccc} & & 1 & & 2 & & 2 \\ & & \circ & \cdots & \circ & => & \circ \\ & & & & | & & \\ & & & & \circ & & \\ & & & & -\rho & & \end{array}$$

where we denote  $-\rho = \alpha_1 + 2\alpha_2 + 2\alpha_3$ .

For the root system  $\Delta$  of  $B_3 = so(7)$  type and  $\dim B_3 = 21$  with simple roots  $\{\alpha_1, \alpha_2, \alpha_3\}$ , it is well known that

$\Delta = \{\alpha_1, \alpha_2, \alpha_3, \alpha_1 + \alpha_2, \alpha_2 + \alpha_3, \alpha_1 + \alpha_2 + \alpha_3, \alpha_2 + 2\alpha_3, \alpha_1 + \alpha_2 + 2\alpha_3, \alpha_1 + 2\alpha_2 + 2\alpha_3\}$ , and

$$B_3 = \mathfrak{g} = \mathfrak{h} \oplus \sum_{\alpha \in \Delta} (\mathfrak{g}_\alpha \oplus \mathfrak{g}_{-\alpha}),$$

where  $\mathfrak{h}$  is the Cartan subalgebra of  $B_3$ .

### 5.1 The case of a JTS

First, we study the case of  $g_{-1} = U = Mat(1, 5; \Phi)$ . (Hereafter, we assume  $\Phi = \mathbb{C}$  complex number field.)

In this case,  $g_{-1}$  is a JTS with respect to the product

$$\{xyz\} = x {}^t yz + y {}^t zx - z {}^t xy,$$

where  ${}^t x$  denotes the transpose matrix of  $x$ .

By straightforward calculations, the standard embedding Lie algebra  $L(U) = \mathfrak{g}$  can be shown to be a 3-graded  $B_3$ -type Lie algebra with  $g_{-1} \oplus g_0 \oplus g_1$ . Thus, we have

$$\begin{aligned} g_0 &= DerU \oplus Anti - DerU \\ &= B_2 \oplus \Phi H, \text{ where } H := \begin{pmatrix} Id & 0 \\ 0 & -Id \end{pmatrix} \\ Der(g_{-1} \oplus g_1) &\cong \{\circ \cdots \circ \implies \circ\} = B_3, \quad (\circ \text{ omitted}). \end{aligned}$$

$$\begin{array}{ccccc} 1 & & 2 & & 2 \\ \circ & \cdots & \circ & \implies & \circ \\ & & | & & \\ & & \circ & & \\ & & -\rho & & \end{array}$$

Furthermore, we obtain

$$\begin{aligned} DerU &= \{L(x, y) - L(y, x)\}_{span} = B_2, \\ Anti - DerU &= \{L(x, y) + L(y, x)\}_{span} = \Phi H, \\ g_0 &= \left\{ \begin{pmatrix} L(x, y) & 0 \\ 0 & -L(y, x) \end{pmatrix} \right\}_{span} = \{S(x, y) + A(x, y)\}_{span}, \end{aligned}$$

where

$$S(x, y) = L(x, y) - L(y, x), A(x, y) = L(x, y) + L(y, x).$$

Here,  $g_{-1}$  corresponds to the root system

$$\{\alpha_1, \alpha_1 + \alpha_2, \alpha_1 + \alpha_2 + \alpha_3, \alpha_1 + \alpha_2 + 2\alpha_3, \alpha_1 + 2\alpha_2 + 2\alpha_3\}.$$

## 5.2 The case of a balanced GJTS

Second, we study the case of  $g_{-1} = U = \text{Mat}(2, 3; \Phi)$ .

In this case,  $g_{-1}$  is a balanced GJTS of 2nd order w.r.t. the product

$$\{xyz\} := z {}^t yx + x {}^t yz - zJ_3 {}^t xyJ_3,$$

where

$$J_3 = \begin{pmatrix} & & 1 \\ & 1 & \\ 1 & & \end{pmatrix}.$$

By straightforward calculations, it can be shown that  $L(U) = g$  is a 5-graded  $B_3$ -type Lie algebra with  $g_{-2} \oplus \cdots \oplus g_2$  and  $\dim g_{-2} = 1$ . Thus, we have

$$g_0 = \text{Der}U \oplus \text{Anti} - \text{Der}U = A_1 \oplus A_1 \oplus \Phi H, \text{ where } H := \begin{pmatrix} \text{Id} & 0 \\ 0 & -\text{Id} \end{pmatrix}$$

$$\text{Der}(g_1 \oplus g_1) = g_{-2} \oplus g_0 \oplus g_2 = A_1 \oplus A_1 \oplus A_1 (\odot \text{ omitted}) \cong \text{Der}T(U).$$

$$\begin{array}{ccccc} 1 & & 2 & & 2 \\ \circ & \cdots & \odot & => & \circ \\ & & | & & \\ & & \circ & & \\ & & -\rho & & \end{array}$$

Furthermore, we obtain

$$g_{-2} = \{K(x, y)\}_{\text{span}} = \Phi \text{Id} \cdots \text{ this means one dimensional,}$$

i.e., balanced. This  $g_{-1}$  corresponds to the root system

$$\{\alpha_2, \alpha_1 + \alpha_2, \alpha_1 + \alpha_2 + \alpha_3, \alpha_2 + \alpha_3, \alpha_2 + 2\alpha_3, \alpha_1 + \alpha_2 + 2\alpha_3\},$$

$g_{-2}$  corresponds to the highest root

$$\{\alpha_1 + 2\alpha_2 + 2\alpha_3\},$$

and  $g/(g_{-2} \oplus g_0 \oplus g_2) \cong T(= g_{-1} \oplus g_1)$  is the tangent space of a quaternion symmetric space of dimension 12, since  $T$  is a Lie triple system associated with  $g_{-1}$  (e.g., [1]).

**5.3 The case of a GJTS of 2nd order**

Third, we study the case of  $g_{-1} = U = Mat(1, 3; \Phi)$ .

In this case,  $g_{-1}$  is a GJTS of 2nd order with respect to the product

$$\{xyz\} = x {}^t yz + z {}^t yx - y {}^t xz.$$

By straightforward calculations, it can be shown that  $L(U)$  is a 5-graded  $B_3$ -type Lie algebra with  $g_{-2} \oplus \dots \oplus g_2$  and  $dim g_{-2} = 3$ ,

$$g_0 = DerU \oplus Anti - DerU = A_2 \oplus \Phi H, \text{ where } H := \begin{pmatrix} Id & 0 \\ 0 & -Id \end{pmatrix}$$

$$Der(g_{-1} \oplus g_1) = g_{-2} \oplus g_0 \oplus g_2 = A_3(\odot \text{ omitted}) \cong DerT(U).$$

$$\begin{array}{cccc} 1 & & 2 & & 2 \\ \circ & \dots & \circ & => & \odot \\ & & | & & \\ & & \circ & & \\ & & -\rho & & \end{array}$$

Furthermore, we obtain

$$g_{-2} = \{K(x, y)\}_{span} = Alt(3, 3; \Phi).$$

That is, the triple system  $g_{-1}$  (resp.  $g_{-2}$ ) corresponds to the root system

$$\{\alpha_3, \alpha_2 + \alpha_3, \alpha_1 + \alpha_2 + \alpha_3\} \text{ (resp. } \{\alpha_2 + 2\alpha_3, \alpha_1 + \alpha_2 + 2\alpha_3, \alpha_1 + 2\alpha_2 + 2\alpha_3\}),$$

implying that

$$\underbrace{\circ \dots \circ \circ \dots \circ}_{\dots} \implies \odot \text{ (}\odot \text{ omitted) and}$$

$$g_0 = A_2 \oplus \Phi H.$$

**Remark.** Following [9], for the case of a GJT of 2nd order, note that  $g_{-2}(\cong \mathbf{k})$  has the structure of the JTS associated with a GJTS of 2nd order.

**5.4 The case of a derivation induced from a JTS**

Finally, we study the case of  $g_{-1} = U = Mat(1, 7; \Phi)$ .

In this case,  $g_{-1}$  is a JTS with respect to the product

$$\{xyz\} = x {}^t yz + y {}^t zx - z {}^t xy.$$

For this case, we obtain

$$DerU = \{L(x, y) - L(y, x)\}_{span} = Alt(7, 7; \Phi) \cong B_3,$$

$$Anti - DerU \cong \Phi H, \text{ which is one dimensional.}$$

The standard embedding Lie algebra is a 3-graded  $B_4$ -type Lie algebra with  $g_{-1} \oplus g_0 \oplus g_1$ .

Furthermore, we have

$$\odot \cdots \odot \underbrace{\cdots \odot \cdots \odot}_{\implies} \odot \quad (\odot \text{ omitted})$$

$$g_0 = B_3 \oplus \Phi H.$$

This case is obtained from  $DerU$  such that  $U = Mat(1, 7; \Phi)$  with the JTS structure.

**Remark.** In the above constructions, note that there exist four different constructions for  $B_3$ -type Lie algebras. These results may be applicable to mathematical physics, for example, quark theory and gravity theory.

## 6 Concluding Remarks

To briefly summarize our study, we note the following.

### 6.1 Geometrical viewpoint

The inner structures of triple systems are closely related to the characterization of root systems of Lie algebras. In particular, we have the following:

(i) There exists a correspondence between simple balanced  $(-1,1)$ -Freudenthal-Kantor triple systems and quaternionic Riemannian symmetric spaces ([1],[12],[17]). That is, there exists a correspondence between simple balanced GJTS of 2nd order and quaternionic Riemannian symmetric spaces.

(ii) There exists a relationship between Lie triple systems and totally geodesic manifolds ([5],[21]).

(iii) There exists a relationship between symmetric domains and positive definite Hermitian Jordan triple systems ([24]).

Thus, triple systems appear to be a useful tool and concept for characterizing geometrical phenomena.

### 6.2 Another viewpoint

For the theory of Peirce decompositions, we refer the reader to ([12],[14],[15]), and for the mathematical physics, we refer the reader to ([6],[19],[22]).

For  $sl_2$  subalgebras, let  $U(\varepsilon, \delta)$  be an  $(\varepsilon, \delta)$ -F-K.t.s. and  $L(\varepsilon, \delta)$  be the standard embedding Lie algebra or superalgebra. If  $P \in \{K(a, b)\}_{span}$  and  $P^2 = \mu Id$ ,  $\sqrt{\mu} \in \Phi$ , then we have the Lie subalgebras:

$$sl_2(\Phi) \leq D(T(\varepsilon, \delta)) \leq L(\varepsilon, \delta) \text{ and } sl_2(\Phi) := \{H, E, F\}_\Phi,$$

where  $E = \frac{1}{\sqrt{\mu}} \begin{pmatrix} 0 & P \\ 0 & 0 \end{pmatrix}$ ,  $F = \frac{1}{\sqrt{\mu}} \begin{pmatrix} 0 & 0 \\ P & 0 \end{pmatrix}$ ,  $H := [E, F] = \begin{pmatrix} Id & 0 \\ 0 & -Id \end{pmatrix}$ .

This means that there is a generalization of  $sl_2(\Phi)$  loop algebra in  $L(\varepsilon, \delta)$ , which we will discuss in a future paper.

### 6.3 7-graded Lie algebra and superalgebra

Let  $g = \sum_{i=-3}^3 g_i$  be a 7 graded Lie algebra or superalgebra such that  $[g_i, g_j] \subseteq g_{i+j}$  if  $|i+j| \leq 3$  and  $[g_i, g_j] = 0$  if  $|i+j| \geq 4$ .

If we define

$$A = g_{-2} \oplus g_{-1} \oplus g_1 \oplus g_2,$$

$$\begin{aligned} D(A, A) &= g_{-3} \oplus g_0 \oplus g_3 \\ &= ad([g_{-2}, g_{-1}] + [g_{-1} + g_{-2}] + [g_{-2}, g_2] + [g_{-1}, g_1] + \\ &\quad + [g_2, g_1] + [g_1, g_2] + [g_2, g_{-2}] + [g_1, g_{-1}]), \end{aligned}$$

where we denote  $ad\ xy = [x, y]$ , then  $A$  is a generalized structurable algebra with respect to the product

$$X \circ Y = [x_{-1}, y_{-1}] + [x_{-2}, y_1] + [x_{-1}, y_2] + [x_1, y_1] + [x_2, y_{-1}] + [x_1, y_{-2}]$$

where  $X = x_{-2} + x_{-1} + x_1 + x_2$  and  $Y = y_{-2} + y_{-1} + y_1 + y_2$ , and the derivation  $D(X, Y)$ . That is, we have

$$D(X \circ Y, Z) + D(Y \circ Z, X) + D(Z \circ X, Y) = 0.$$

Finally, we emphasize that many mathematical and physical subjects involving geometric phenomena may be characterized by applying concept of triple systems.

## References

- [1] W.Bertram, *Complex and quaternionic structures on symmetric spaces -correspondence with Freudenthal-Kantor triple systems*, Sophia Kokyuroku in Math., **45**, Theory of Lie Groups and Manifolds, ed. R.Miyaoka and T.Tamaru, 57–76, 2002.
- [2] A.Elduque, N.Kamiya and S.Okubo, *Simple (-1,-1) balanced Freudenthal-Kantor triple systems*, Glasgow Math. J., **45**, 353–372, 2003.
- [3] A.Elduque, N.Kamiya and S.Okubo, *(-1,-1) balanced Freudenthal-Kantor triple systems and noncommutative Jordan algebras*, J. Alg., **294**, 19–40, 2005.

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- [4] H.Freudenthal and H.de Vries, *Linear Lie Groups*, Acad. Press, New York, 1969.
  - [5] S.Helgason, *Differential Geometry, Lie Groups, and Symmetric Spaces*, Acad. Press, New York, 1978.
  - [6] R.Iordanescu, *Jordan structures in Analysis, Geometry and Physics*, Editura Acad. Romane, Bucuresti, 2009.
  - [7] N.Jacobson, *Structures and Representations of Jordan Algebras*, Amer. Math. Soc. Colloq. **39**, Amer. Math. Soc., Providence, RI, 1968.
  - [8] N.Kamiya, *A structure theory of Freudenthal-Kantor triple systems*, J. Alg., **110**, 108–123, 1987.
  - [9] N.Kamiya, *A structure theory of Freudenthal-Kantor triple systems II*, Comm. Math. Univ. Sancti., **38**, 41–60, 1989.
  - [10] N.Kamiya, *On radicals of triple systems*, Groups, Rings, Lie and Hopf Algebras, (St. John's 2001), 75–83, Math. Appl., 555, Kluwer Acad. Publ., Dordrecht, 2003.
  - [11] N.Kamiya, *On a realization of the exceptional simple graded Lie algebras of second kind and Freudenthal-Kantor triple systems*, Polish Academy of Sciences Math. **40**, no.1, 55–65, 1998.
  - [12] N.Kamiya, *Examples of Peirce decomposition of generalized Jordan triple systems of second order – Balanced cases–*, Contemporary Mathematics, **391**, A.M.S., 157–166, 2005.
  - [13] N.Kamiya, *On a generalization of structurable algebras*, Algebras, Groups and Geometries, **9**, 35–47, 1992.
  - [14] I.L.Kantor and N.Kamiya, *A Peirce decomposition for generalized Jordan triple systems of second order*, Comm. Alg., **31**, no.12, 5875–5913, 2003.
  - [15] N.Kamiya and D.Mondoc, *Examples of Peirce decomposition of Kantor triple systems*, Algebras, Groups and Geometries, 24, 325–348, 2007.
  - [16] N.Kamiya and D.Mondoc, *A new class of nonassociative algebras with involution*, Proc. Japan Acad. Ser.A, 84, no.5, 68–72, 2008.
  - [17] N.Kamiya and S.Okubo, *On  $\delta$ -Lie Supertriple Systems Associated with  $(\varepsilon, \delta)$ -Freudenthal-Kantor Supertriple Systems*, Proc. Edinburgh Math. Soc., **43**, 243–260, 2000.

- [18] N.Kamiya and S.Okubo, *Construction of Lie superalgebras  $D(2, 1; \alpha)$ ,  $G(3)$  and  $F(4)$  from some triple systems*, Proc. Edinburgh Math. Soc., **46**, 87–98, 2003.
- [19] N.Kamiya and S.Okubo, *On generalized Freudenthal-Kantor triple systems and Yang-Baxter equations*, Proc. XXIV International Coll. Group Theoretical Methods in Physics, Inst. Physics Conf. Ser., **173**, 815–818, 2003.
- [20] N.Kamiya and S.Okubo, *On composition, quadratic and some triple systems*, Lecture Notes in Pure and App. Math., **246**, Taylor (CRC), 205–231, 2006.
- [21] O.Loos, *Symmetric Spaces*, Benjamin, London, 1969.
- [22] S.Okubo, *Introduction to Octonion and other Non-associative Algebras in Physics*, Cambridge Univ. Press, Cambridge, 1995.
- [23] S.Okubo and N.Kamiya, *Jordan-Lie superalgebras and Jordan-Lie triple systems*, J. Alg., **198**, 388–411, 1997.
- [24] I.Satake, *Algebraic Structures of Symmetric Domains*, Princeton Univ. Press, Tokyo, 1980.
- [25] M.Scheunert, *The Theory of Lie Superalgebras*, Lecture Notes in Math. **716**, Springer, New York, 1979.

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